Stability of standing waves for NLS equations with a delta function potential

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We consider the following nonlinear Schrödinger equation with a delta function potential:

$$i\partial_t u = -\partial_x^2 u + \gamma \delta(x)u + \alpha |u|^{p-1}u, \quad (t,x) \in \mathbb{R} \times \mathbb{R},$$
(1)

where $\gamma, \alpha \in \mathbb{R}, 1 , and <math>\delta(x)$ is the delta function at x = 0. The formal expression $-\partial_x^2 + \gamma \delta(x)$ in (1) is formulated as a linear operator A_{γ} or H_{γ} associated with a quadratic form a_{γ} on $H^1(\mathbb{R})$:

$$a_{\gamma}(u,v) = \int_{\mathbb{R}} \partial_x u(x) \overline{\partial_x v(x)} \, dx + \gamma u(0) \overline{v(0)}, \quad u,v \in H^1(\mathbb{R})$$

Then, there exists a unique bounded linear operator $A_{\gamma} : H^1(\mathbb{R}) \to H^{-1}(\mathbb{R})$ such that

 $\langle A_{\gamma}u, v \rangle = \operatorname{Re} a_{\gamma}(u, v), \quad u, v \in H^1(\mathbb{R}).$

Moreover, we define a linear operator H_{γ} in $L^2(\mathbb{R})$ by $H_{\gamma}v = -\partial_x^2 v$ for $v \in D(H_{\gamma})$ with the domain

$$D(H_{\gamma}) = \{ v \in H^2(\mathbb{R} \setminus \{0\}) \cap H^1(\mathbb{R}) : \partial_x v(+0) - \partial_x v(-0) = \gamma v(0) \}.$$

Then, H_{γ} is a self-adjoint operator in $L^2(\mathbb{R})$, and satisfies

$$(H_{\gamma}u, v)_{L^2} = a_{\gamma}(u, v), \quad u, v \in D(H_{\gamma}).$$

The following spectral properties of H_{γ} are known: $\sigma_{\rm ess}(H_{\gamma}) = \sigma_{\rm ac}(H_{\gamma}) = [0,\infty), \ \sigma_{\rm sc}(H_{\gamma}) = \emptyset$. If $\gamma \geq 0, \ \sigma_{\rm p}(H_{\gamma}) = \emptyset$. If $\gamma < 0, \ \sigma_{\rm p}(H_{\gamma}) = \{-\gamma^2/4\}$ with its positive normalized eigenfuction $(|\gamma|/2)^{1/2}e^{-|\gamma||x|/2}$ (see [1, Chapter I.3] for details).

In this talk, we study the structure and the orbital stability of standing wave solutions $e^{i\omega t}\varphi_{\omega}(x)$ for (1), where $\omega \in \mathbb{R}$ is a parameter, and $\varphi_{\omega} \in H^1(\mathbb{R})$ is a positive solution of the stationary problem:

$$A_{\gamma}\varphi + \omega\varphi + \alpha |\varphi|^{p-1}\varphi = 0 \quad \text{in } H^{-1}(\mathbb{R}).$$
⁽²⁾

The local well-posedness of the Cauchy problem for (1) in the energy space $H^1(\mathbb{R})$ follows from an abstract result in Cazenave [2]. Moreover, there is conservation of charge and energy, i.e.,

$$||u(t)||_{L^2} = ||u_0||_{L^2}, \quad E(u(t)) = E(u_0), \quad \forall t \in [0, T^*),$$

where u(t) is the solution of (1) with $u(0) = u_0 \in H^1(\mathbb{R}), T^* = T^*(u_0) \in (0, \infty]$ is the maximal existence time of u(t), and E is defined by

$$E(v) = \frac{1}{2}a_{\gamma}(v,v) + \frac{\alpha}{p+1} \|v\|_{L^{p+1}}^{p+1}$$
$$= \frac{1}{2} \|\partial_x v\|_{L^2}^2 + \frac{\gamma}{2} |v(0)|^2 + \frac{\alpha}{p+1} \|v\|_{L^{p+1}}^{p+1}$$

(see Theorem 3.7.1 and Corollary 3.3.11 in [2]).

For the stability of standing waves for (1), the case where $\gamma < 0$ and $\alpha < 0$ was first studied by Goodman, Holmes and Weinstein [7] for the special case p = 3, and then by Fukuizumi, Ohta and Ozawa [6] for general case 1 .

Theorem 1 ([6]) Let $\gamma < 0$, $\alpha = -1$ and $1 . If <math>\omega > \gamma^2/4$, the stationary problem (2) has a unique positive solution $\varphi_{\omega} \in H^1(\mathbb{R})$ given by

$$\varphi_{\omega}(x) = \left(\frac{(p+1)\omega}{2}\right)^{1/(p-1)} \left\{ \cosh\left(\frac{(p-1)\sqrt{\omega}}{2}|x| + b(\omega)\right) \right\}^{-2/(p-1)}$$
(3)

for $x \in \mathbb{R}$, where $b(\omega) = \tanh^{-1}(-\frac{\gamma}{2\sqrt{\omega}})$. If $1 , the standing wave solution <math>e^{i\omega t}\varphi_{\omega}$ of (1) is stable for any $\omega \in (\gamma^2/4, \infty)$. If p > 5, there exists $\omega^* = \omega^*(\gamma, p) \in (\gamma^2/4, \infty)$ such that $e^{i\omega t}\varphi_{\omega}$ is stable for any $\omega \in (\gamma^2/4, \omega^*)$, and is unstable for any $\omega \in (\omega^*, \infty)$.

Remark that for the case where $\gamma = 0$ and $\alpha < 0$, the standing wave solution $e^{i\omega t}\varphi_{\omega}$ is stable for any $\omega \in (0, \infty)$ if 1 , and is unstable for $any <math>\omega \in (0, \infty)$ if $p \ge 5$. Fukuizumi and Jeanjean [5] studies the case where $\gamma > 0$ and $\alpha < 0$. Note that the expression (3) of φ_{ω} is valid for the case where $\gamma > 0$ and $\alpha = -1$ (see also [4] for the expression of φ_{ω}). The stability of $e^{i\omega t}\varphi_{\omega}$ is determined by the sign of the derivative of the function

$$\omega \mapsto \|\varphi_{\omega}\|_{L^{2}}^{2} = C_{p} \omega^{\frac{5-p}{2(p-1)}} \int_{b(\omega)}^{\infty} (\cosh y)^{-4/(p-1)} dy$$

where C_p is a positive constant depending only on p. Moreover, the fact that φ_{ω} is a minimizer of the minimization problem

$$\inf\{S_{\omega}(v): v \in H^1(\mathbb{R}) \setminus \{0\}, \ I_{\omega}(v) = 0\}$$

plays an important role in the proof of Theorem 1. Here, we put

$$S_{\omega}(v) = \frac{1}{2} \|\partial_x v\|_{L^2}^2 + \frac{\omega}{2} \|v\|_{L^2}^2 + \frac{\gamma}{2} |v(0)|^2 - \frac{1}{p+1} \|v\|_{L^{p+1}}^{p+1},$$

$$I_{\omega}(v) = \|\partial_x v\|_{L^2}^2 + \omega \|v\|_{L^2}^2 + \gamma |v(0)|^2 - \|v\|_{L^{p+1}}^{p+1}.$$

Next, we consider the case where $\gamma < 0$ and $\alpha > 0$ (attractive potential and repulsive nonlinearity).

Theorem 2 Let $\gamma < 0$, $\alpha = 1$ and $1 . If <math>0 < \omega < \gamma^2/4$, the stationary problem (2) has a unique positive solution $\varphi_{\omega} \in H^1(\mathbb{R})$ given by

$$\varphi_{\omega}(x) = \left(\frac{(p+1)\omega}{2}\right)^{1/(p-1)} \left\{ \sinh\left(\frac{(p-1)\sqrt{\omega}}{2}|x| + c(\omega)\right) \right\}^{-2/(p-1)}$$

for $x \in \mathbb{R}$, where $c(\omega) = \tanh^{-1}(2\sqrt{\omega}/|\gamma|)$. The standing wave solution $e^{i\omega t}\varphi_{\omega}$ of (1) is orbitally stable.

Theorem 3 Let $\gamma < 0$, $\alpha = 1$ and $\omega = 0$. If $1 , the stationary problem (2) has a unique positive solution <math>\varphi_0 \in H^1(\mathbb{R})$ given by

$$\varphi_0(x) = \left(\frac{2(p+1)\gamma^2}{\{4+(p-1)|\gamma||x|\}^2}\right)^{1/(p-1)}$$

for $x \in \mathbb{R}$. The stationary solution φ_0 of (1) is orbitally stable. If $p \ge 5$, the stationary problem (2) has no nontrivial solutions in $H^1(\mathbb{R})$.

The proof of Theorem 2 is based on the fact that φ_{ω} is characterized by a minimizer of the minimization problem

$$\inf\{S_{\omega}(v): v \in H^1(\mathbb{R})\},\$$

and the conservation of energy and charge (cf. Cazenave and Lions [3]). Theorem 3 is proved in a similar way.

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